

Suggested Solutions to Midterm Test for MATH4220

March 10, 2016

1. (20 points)

(a) (8 points) Find the general solutions to

$$u_x - 2u_y + 2u = 0$$

(b) (12 points) Solve the problem:

$$\begin{cases} y\partial_x u + 3x^2 y \partial_y u = 0 \\ u(x=0, y) = y^2 \end{cases}$$

Find the region in the xy -plane so that the solution is uniquely determined.

Solution:

(a) **Method 1: Coordinate Method:**

Change variables to

$$x' = x - 2y, \quad y' = -2x - y$$

Hence $u_x - 2u_y + 2u = 5u_{x'} + 2u = 0$. Thus the solution is $u(x', y') = f(y')e^{-\frac{2}{5}x'}$, with f an arbitrary function of one variable. Therefore, the general solutions are

$$u(x, y) = f(-2x - y)e^{-\frac{2}{5}(x-2y)}$$

where f is an arbitrary function.

Method 2: Geometric Method

The corresponding characteristic curves are

$$\frac{dx}{1} = \frac{dy}{-2}$$

that is, $y = -2x + C$ where C is an arbitrary constant. Then

$$\frac{d}{dx}u(x, -2x + C) = u_x(x, -2x + C) - 2u_y(x, -2x + C) = -2u(x, -2x + C)$$

Hence $u(x, -2x + C) = f(C)e^{-2x}$, where f is an arbitrary function. Therefore,

$$u(x, y) = f(2x + y)e^{-2x}$$

where f is an arbitrary function.

(b) The characteristic curves are

$$\frac{dy}{3x^2 y} = \frac{dx}{y}$$

that is, $y = x^3 + C$ where C is an arbitrary constant. Then

$$\frac{d}{dx}u(x, x^3 + C) = u_x + 3x^2 u_y = 0$$

Hence $u(x, x^3 + C) = f(C)$ where f is an arbitrary function. Thus

$$u(x, y) = f(y - x^3)$$

Besides, the auxiliary condition gives that $y^2 = u(x = 0, y) = f(y)$. Hence, the solution is

$$u(x, y) = (y - x^3)^2$$

Note that when $y = 0$ the equation vanishes, thus the characteristic curves break down when $y = 0$, therefore the solution is uniquely determined on $\{(x, y) : y > 0, y > x^3\} \cup \{(x, y) : y < 0, y < x^3\} \cup \{(0, 0)\}$. (Remark: if the solution is **continuous**, then u is uniquely determined on the whole plane by the continuity of u).

2. (30 points)

(a) (5 points) State the definition of a well-posed PDE problem.

(b) (5 points) Is the following problem well-posed? Why?

$$\begin{cases} \partial_x^2 u + \partial_y^2 u = 0, & x^2 + y^2 < 1 \\ \frac{\partial u}{\partial \vec{n}}(x, y) = 0, & x^2 + y^2 = 1, \vec{n} \text{ is the unit outnorm of } x^2 + y^2 = 1 \end{cases}$$

(c) (10 points) Verifying that $u_n(x, t) = \frac{1}{n} \sin nx e^{-n^2 t}$ solves the following problem

$$\begin{cases} \partial_t u = \partial_x^2 u, & 0 < x < \pi, -\infty < t < +\infty \\ u(0, t) = u(\pi, t) = 0, & -\infty < t < \infty \\ u(x, t = 0) = \frac{1}{n} \sin nx, & 0 \leq x \leq \pi \end{cases}$$

for all positive integer n .

How does the energy change when $t \rightarrow \pm\infty$?

(d) (10 points) Is the following problem

$$\begin{cases} \partial_t u = \partial_x^2 u, & 0 < x < \pi, t < 0 \\ u(0, t) = u(\pi, t) = 0, & t < 0 \\ u(x, t = 0) = 0, & 0 < x < \pi \end{cases}$$

well-posed? Why?

Solution:

(a) A PDE problem is said to be well-posed if the following three properties are satisfied:

Existence: There exists at least one solution $u(x, t)$ satisfying all these conditions.

Uniqueness: There is at most one solution.

Stability: The unique solution $u(x, t)$ depends in a stable manner on the data of the problem. This means that if the data are changed a little, the corresponding solution changes only a little.

(b) No.

Let $u(x, t) = C$ where C is an arbitrary constant. Then $\partial_x u = \partial_y u = \partial_x^2 u = \partial_y^2 u = 0$, hence

$$\partial_x^2 u + \partial_y^2 u = 0, x^2 + y^2 < 1$$

$$\frac{\partial u}{\partial \vec{n}} = (\partial_x u, \partial_y u) \cdot (x, y) = 0, x^2 + y^2 = 1$$

Therefore, any constant is the solution of the problem. Hence the solution exists but is not unique.

(c) After a little simple computations, we have for all positive interger n

$$\partial_t u_n(x, t) = -n \sin(nx) e^{-n^2 t}$$

$$\partial_x u_n(x, t) = \cos(nx) e^{-n^2 t}$$

$$\partial_x^2 u_n(x, t) = -n \sin(nx) e^{-n^2 t}$$

then $\partial_t u_n = -n \sin(nx) e^{-n^2 t} = \partial_x^2 u_n$, $0 < x < \pi$, $-\infty < t < \infty$. And

$$u_n(0, t) = 0, \quad -\infty < t < \infty$$

$$u_n(\pi, t) = 0, \quad -\infty < t < \infty$$

$$u_n(x, t = 0) = \frac{1}{n} \sin(nx), \quad 0 < x < \pi$$

hence u_n is indeed the solution of the problem.

The energy is

$$E = \frac{1}{2} \int_0^\pi |u_n(x, t)|^2 dx = \frac{1}{2n^2} e^{-2n^2 t} \int_0^\pi \sin^2(nx) dx = \frac{\pi}{4n^2} e^{-2n^2 t}$$

hence $E(t) \rightarrow 0$ as $t \rightarrow +\infty$ and $E(t) \rightarrow +\infty$ as $t \rightarrow -\infty$.

(d) No.

On one hand, $u = 0$ is a solution of

$$\begin{cases} \partial_t u = \partial_x^2 u, & 0 < x < \pi, \quad t < 0 \\ u(0, t) = u(\pi, t) = 0, & t < 0 \\ u(x, t = 0) = 0, & 0 < x < \pi \end{cases}$$

On the other hand, $u_n(x, t) = \frac{1}{n} \sin nx e^{-n^2 t}$ solves the following problem

$$\begin{cases} \partial_t u = \partial_x^2 u, & 0 < x < \pi, \quad t < 0 \\ u(0, t) = u(\pi, t) = 0, & t < 0 \\ u(x, t = 0) = \frac{1}{n} \sin nx, & 0 \leq x \leq \pi \end{cases}$$

for all positive integer n by above problem (c).

Note that

$$\int_0^\pi \left| \frac{1}{n} \sin nx - 0 \right|^2 dx = \frac{\pi}{2n^2} \rightarrow 0 \quad \text{as } n \rightarrow +\infty$$

However when $t < 0$,

$$\int_0^\pi \left| \frac{1}{n} \sin(nx) e^{-n^2 t} - 0 \right|^2 dx = \frac{\pi}{2n^2} e^{-2n^2 t} \rightarrow +\infty \quad \text{as } n \rightarrow \infty$$

Hence, when the data $u(x, t = 0)$ changes a little in the sense of L^2 -norm, the difference of the solutions in L^2 -norm tends to infinity. This violates the stability in the sense of L^2 -norm, therefore, it is not well-posed.

Remark: consider the stability in uniform sense.

On one hand, $\max_{0 < x < \pi} \left| \frac{1}{n} \sin(nx) - 0 \right| \rightarrow 0$ as $n \rightarrow \infty$

On the other hand, when $t < 0$, $\max_{0 < x < \pi} \left| \frac{1}{n} \sin(nx) e^{-n^2 t} - 0 \right| \rightarrow +\infty$ as $n \rightarrow \infty$

This violates the stability in the uniform sense.

3. (10 points) Is there a maximum principle for the Cauchy problem for the 1-dimensional wave equation? Explain why?

Solution:No.

Consider the following Cauchy problem:

$$\begin{cases} \partial_t^2 u - \partial_x^2 u = 0, & -\infty < x < +\infty, \quad t > 0 \\ u(x, t = 0) = 0, \quad \partial_t u(x, t = 0) = \sin x, & -\infty < x < +\infty \end{cases}$$

And the unique solution is given by d'Alembert formula:

$$u(x, t) = \frac{1}{2} \cos(x+t) - \cos(x-t) = -\sin x \sin t, \quad -\infty < x < \infty, t > 0$$

Then $u(x, t)$ attains its maximum 1 only at the interior points $(\frac{\pi}{2} \pm 2n\pi, \frac{3\pi}{2} + 2n\pi)$ or $(\frac{3\pi}{2} \pm 2n\pi, \frac{\pi}{2} + 2n\pi)$ for $n = 0, 1, 2, \dots$. However, $u(x, t) = 0$ on the boundary $\{(x, t) : t = 0\}$. Therefore there is no maximum principle for the Cauchy problem for the 1-dimensional wave equation.

Remark: The key is to find a counterexample.

4. (10 points)

- (a) (5 points) What is the type of the equation

$$\partial_t^2 u + \partial_{xt}^2 u - 2\partial_x^2 u = 0 ?$$

- (b) (5 points) Solve the Cauchy problem

$$\begin{cases} \partial_t^2 u - 2\partial_x^2 u = 0, & -\infty < x < +\infty, \quad -\infty < t < +\infty \\ u(x, t = 0) = x^2, \quad \partial_t u(x, t = 0) = \sin x, & -\infty < x < +\infty \end{cases}$$

Solution:

- (a) Since $a_{11} = 1, a_{12} = \frac{1}{2}, a_{22} = -2$, then $a_{12}^2 - a_{11}a_{22} = \frac{9}{4} > 0$, hence it is hyperbolic.
 (b) The solution is given by d'Alembert Formula directly:

$$u(x, t) = \frac{1}{2}[\phi(x+ct) + \phi(x-ct)] + \frac{1}{2c} \int_{x-ct}^{x+ct} \psi y dy$$

Here $c = \sqrt{2}$, $\phi(x) = x^2$ and $\psi(x) = \sin x$. Hence

$$u(x, t) = \frac{1}{2}[(x+ct)^2 + (x-ct)^2] + \frac{1}{2\sqrt{2}} \int_{x-ct}^{x+ct} \sin y dy = x^2 + 2t^2 + \frac{1}{\sqrt{2}} \sin x \sin \sqrt{2}t$$

5. (20 points) Consider the Cauchy problem

$$\begin{cases} \partial_t u = \partial_x^2 u, & -\infty < x < +\infty, \quad t > 0 \\ u(x, t = 0) = \phi(x) & -\infty < x < +\infty \end{cases}$$

- (a) (10 points) Show that any finite energy solution to the Cauchy problem is unique by the energy method.
 (b) (10 points) Find the solution with $\phi(x)$ given by

$$\phi(x) = \begin{cases} 1, & \text{if } x > 0 \\ 0, & \text{if } x < 0 \end{cases}$$

Solution:

- (a) Suppose u_1 and u_2 are two finite energy solution to Cauchy problem. Let $v(x, t) = u_1(x, t) - u_2(x, t)$, then $v(x, t)$ satisfies the following problem:

$$\begin{cases} \partial_t v = \partial_x^2 v, & -\infty < x < +\infty, \quad t > 0 \\ v(x, t = 0) = 0 & -\infty < x < +\infty \end{cases}$$

Multiplying the both sides of $\partial_t v = \partial_x^2 v$ by v and taking intergral from $-\infty$ to ∞ with respect to x , then we have

$$\int_{-\infty}^{\infty} \partial_t v v dx = \int_{-\infty}^{\infty} \partial_x^2 v v dx$$

Then

$$L.H.S = \frac{d}{dt} \int_{-\infty}^{\infty} \frac{1}{2} v^2 dx$$

$$R.H.S = \partial_x v v \Big|_{-\infty}^{\infty} - \int_{-\infty}^{\infty} (\partial_x v)^2 dx$$

Note that u_1 and u_2 are finite energy solutions, that is,

$$\frac{1}{2} \int_{-\infty}^{\infty} |u_1(x, t)|^2 dx < +\infty, \quad \frac{1}{2} \int_{-\infty}^{\infty} |u_2(x, t)|^2 dx < +\infty$$

then

$$\frac{1}{2} \int_{-\infty}^{\infty} |v(x, t)|^2 dx \leq \int_{-\infty}^{\infty} |u_1(x, t)|^2 + |u_2(x, t)|^2 dx < +\infty$$

that is, $v(x, t)$ is a finite energy solution which implies that $v(x, t) \rightarrow 0$ as $x \rightarrow \pm\infty$. Hence

$$R.H.S = - \int_{-\infty}^{\infty} (\partial_x v)^2 dx$$

Then, we have

$$\frac{d}{dt} \int_{-\infty}^{\infty} \frac{1}{2} v^2 dx = - \int_{-\infty}^{\infty} (\partial_x v)^2 dx \leq 0$$

and then for $t > 0$

$$0 \leq \int_{-\infty}^{\infty} \frac{1}{2} v^2(x, t) dx \leq \int_{-\infty}^{\infty} \frac{1}{2} v^2(x, 0) dx = 0$$

By the continuity of v , we have $v(x, t) \equiv 0$, $-\infty < x < \infty, t > 0$. Thus we have shown that $u_1(x, t) \equiv u_2(x, t)$ for $-\infty < x < \infty, t > 0$. Therefore any finite energy solution is unique.

- (b) **Method 1:** Find a solution with the form $u(x, t) = U(\frac{x}{\sqrt{4t}})$. Let $p = \frac{x}{\sqrt{4t}}$, then $U(\frac{x}{\sqrt{4t}})$ satisfies the following equation:

$$U''(p) + 2pU'(p) = 0$$

Hence $U(p) = C_1 + C_2 \int_0^p e^{-s^2} ds$ where C_1, C_2 are arbitrary constants. That is,

$$u(x, t) = U\left(\frac{x}{\sqrt{4t}}\right) = C_1 + C_2 \int_0^{\frac{x}{\sqrt{4t}}} e^{-s^2} ds$$

Now use the initial condition, expressed as a limit as follows.

$$x > 0, \quad 1 = \lim_{t \rightarrow 0^+} U\left(\frac{x}{\sqrt{4t}}\right) = C_1 + C_2 \int_0^{+\infty} e^{-s^2} ds$$

$$x < 0, \quad 0 = \lim_{t \rightarrow 0^+} U\left(\frac{x}{\sqrt{4t}}\right) = C_1 + C_2 \int_0^{-\infty} e^{-s^2} ds$$

Note that $\int_0^{\infty} e^{-s^2} ds = \frac{\sqrt{\pi}}{2}$, we have $C_1 = \frac{1}{2}, C_2 = \frac{1}{\sqrt{\pi}}$. Therefore,

$$u(x, t) = \frac{1}{2} + \frac{1}{\sqrt{\pi}} \int_0^{\frac{x}{\sqrt{4t}}} e^{-s^2} ds$$

Method 2: The solution for the Cauchy problem is given by

$$u(x, t) = \int_{-\infty}^{\infty} S(x - y, t) \phi(y) dy$$

where $S(x, t)$ is the heat kernel, $S(x, t) = \frac{1}{\sqrt{4k\pi t}} e^{-\frac{x^2}{4kt}}$. Here, $k = 1$

$$\phi(x) = \begin{cases} 1, & \text{if } x > 0 \\ 0, & \text{if } x < 0 \end{cases}$$

Then,

$$u(x, t) = \int_0^{\infty} \frac{1}{\sqrt{4\pi t}} e^{-\frac{(x-y)^2}{4t}} dy = \frac{1}{2} + \frac{1}{\sqrt{\pi}} \int_0^{\frac{x}{\sqrt{4t}}} e^{-s^2} ds.$$

6. (10 points) Derive the solution formula for the following initial-boundary value problem

$$\begin{cases} \partial_t u = \partial_x^2 u, & 0 < x < +\infty, \quad t > 0 \\ u(x, t = 0) = \phi(x) & 0 < x < +\infty \\ \partial_x u(x = 0, t) = 0, & t > 0 \end{cases}$$

by the method of reflection.

Solution: Use the reflection method, and first consider the following Cauchy Problem:

$$\begin{cases} \partial_t v = \partial_x^2 v, & 0 < x < +\infty, \quad t > 0 \\ v(x, t = 0) = \phi_{\text{even}}(x) & 0 < x < +\infty \end{cases}$$

where $\phi_{\text{even}}(x)$ is even extension of ϕ which is given by

$$\phi_{\text{even}}(x) = \begin{cases} \phi(x), & \text{if } x > 0 \\ \phi(-x), & \text{if } x < 0 \end{cases}$$

Then the unique solution is given by:

$$v(x, t) = \int_{-\infty}^{\infty} S(x - y, t) \phi_{\text{even}}(y) dy$$

And since $\phi_{\text{even}}(x)$ is even, so is $v(x, t)$ for $t > 0$, which implies

$$\partial_x v(x = 0, t) = 0, t > 0$$

Set $u(x, t) = v(x, t), x > 0$, then $u(x, t)$ is the unique solution of Neumann Problem on the half-line. More precisely, $x > 0, t > 0$

$$\begin{aligned} u(x, t) &= \int_0^{\infty} S(x - y, t) \phi(y) dy + \int_{-\infty}^0 S(x - y, t) \phi(-y) dy \\ &= \int_0^{\infty} S(x - y, t) \phi(y) dy + \int_0^{\infty} S(x + y, t) \phi(y) dy \\ &= \frac{1}{\sqrt{4k\pi t}} \int_0^{\infty} [e^{-\frac{(x-y)^2}{4kt}} + e^{-\frac{(x+y)^2}{4kt}}] \phi(y) dy. \end{aligned}$$